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Energy Extraction Performance Improvement of a Flapping Foil by the Use of Combined Foil

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ABSTRACT

In this study, numerical investigations on the energy extraction performance of a flapping foil device are carried out by using a modified foil shape. The new foil shape is designed by combining the thick leading edge of NACA0012 foil and the thin trailing edge of NACA0006 foil. The numerical simulations are based on the solution of the unsteady and incompressible Navier-Stokes equations that govern the fluid flow around the flapping foil. These equations are resolved in a two-dimensional domain with a dynamic mesh technique using the CFD software ANSYS Fluent 16. A User Define Function (UDF) controls the imposed sinusoidal heaving and pitching motions. First, for a validation study, numerical simulations are performed for a NACA0012 foil undergoing imposed heaving and pitching motions at a low Reynolds number. The obtained results are in good agreement with numerical and experimental data available in the literature. Thereafter, the computations are applied for the new foil shape. The influences of the connecting area location between the leading and trailing segments, the Strouhal number and the effective angle of attack on the energy extraction performance are investigated at low Reynolds number ($Re = 10\,000$). Then, the new foil shape performance was compared to those of both NACA0006 and NACA0012 baseline foils. The results have shown that the proposed foil shape achieves higher performance compared to the baseline NACA foils. Moreover, the energy extraction efficiency was improved by 30.60% compared to NACA0006 and by 17.32% compared to NACA0012. The analysis of the flow field around the flapping foils indicates a change of the vortex structure and the pressure distribution near the trailing edge of the combined foil compared to the baseline foils.

Keywords: Flapping foil; Energy extraction; Power coefficient; Combined foil; CFD.

NOMENCLATURE

c	foil chord length	h_0	nondimensional heaving amplitude
C_D	drag coefficient	$\bar{}$	mean value over one motion cycle
C_L	lift coefficient	$M_z(t)$	instantaneous moment
C_M	moment coefficient	$P(t)$	instantaneous total power extracted
C_{OP}	power coefficient	P_a	total power available in flow
C_{OPh}	power coefficient of heaving motion	$P_h(t)$	instantaneous power extracted by the heaving motion
$C_{OP\theta}$	power coefficient of pitching motion	$P_\theta(t)$	instantaneous power extracted by the pitching motion
CP	pressure coefficient	Re	Reynolds number ($Re = \rho c U_\infty / \mu$)
d	maximum vertical displacement of the trailing edge	St	Strouhal number ($St = 2ch_0f/U_\infty$)
f	flapping frequency	St_{TE}	Strouhal number based on the trailing edge excursion ($St_{TE} = 2df/U_\infty$)
$F_y(t)$	instantaneous vertical force		
$h(t)$	heaving motion		

T	flapping period ($T = 1/f$)	θ_0	nondimensional pitching amplitude
U_∞	free stream velocity	μ	dynamic viscosity
X_P	chordwise position of pitching axis	ρ	fluid density
X_s	connecting area location	\emptyset	phase angle between heaving and pitching motions
α_{eff}	effective angle of attack	χ	feathering parameter
η	energy extraction efficiency	ω	angular frequency ($\omega = 2\pi f$)
$\theta(t)$	pitching motion		

1. INTRODUCTION

The global demand on energy is rapidly increasing as a result of population growth and economic development. Up to date, the world's energy is mainly generated from fossil fuel sources such as oil, coal and gas. Increasing depletion of fossil resources and the rising concern on GHG emissions have led to increased efforts to find sustainable alternative energy sources such as solar, wind and tidal. In this context, harnessing kinetic energy available in the flows is a very interesting prospect. Traditionally, this type of energy is recovered using conventional turbine with rotating blades (horizontal-axis or vertical-axis turbines). Recently, a bio-inspired flapping foil mechanism has offered an alternative strategy to convert the kinetic energy of wind or tidal into mechanical energy and subsequently into electrical energy. At the beginning, the flapping foil mechanism was used as a thrust generator for micro air vehicles (MAV) or autonomous underwater vehicles (AUV). Afterwards, several studies have found that a flapping foil can switch from a propulsive mode (energy consumption) to a power extraction mode (energy harvesting) if the pitch angle is higher to the induced angle of attack (Jones and Platzer 1997). Flapping foils based energy extraction systems are classified into three types of models according to the actuating mechanism: fully-activated systems, semi-activated systems and self-sustained systems (Xiao and Zhu 2014). In fully-activated systems, the foil undergoes imposed heaving and pitching motions. In semi-activated systems, the pitching motion is activated by energy input, while the recovered energy is ensured through the heaving motion caused by the dynamic lifting forces. The net energy extracted is the difference between the energy consumed to activate the pitching motion and the energy recovered by the induced heaving motion. In self-sustained systems, both pitching and heaving motions are generated by the flow-induced instability and therefore do not require actuation device.

Wu (1972) was the first who investigated the energy extraction from flow using flapping foil. His theoretical results showed that a flapping foil in combined heaving and pitching motion can extract energy from an oscillatory flow. Based on this concept, McKinney and DeLaurier (1981) developed the first energy extraction device based on the flapping motion. Their experimental tests in a wind tunnel revealed that the proposed system (called windmill) was competitive compared to other power generation devices. Later, experimental and numerical investigations have been done successively by Davids (1999), Kinsey and Dumas

(2008), Simpson and Triantafyllou (2008), Ashraf *et al.* (2011) and Lu *et al.* (2015). All these works confirmed the possibility of using flapping foil to extract energy. In addition, it was observed that the formation and evolution of the leading edge vortex (LEV) play an important role in the energy extraction process. Their results reported that the best efficiency was always achieved when the reduced frequency ($f^* = fc/U$) was in the range 0.10 – 0.15, the phase difference between the pitching and plunging motions \emptyset was equal or close to 90° and the maximum effective angle of attack α_{ef} was around 30° . At such conditions, a good synchronization between the LEV shedding and the heaving motion was observed.

Afterward, numerous attempts have been made to improve the energy extraction performance of a flapping foil. Some researchers focused on using two flapping foils in tandem. This was the case for (Ashraf *et al.* 2011) that investigated the effects of the phase difference and the distance between the two foils. They found out that in the tandem configuration both averaged power coefficient and efficiency per airfoil were reduced by around 20% compared to a single foil. However, the total efficiency of the optimal tandem configuration increased by 59% compared to a single foil. In a similar study, Kinsey and Dumas (2012) found that optimal tandem configuration has provided high extraction efficiency (up to 64%).

Non-sinusoidal motions were also used to improve the power extraction performance. Ashraf *et al.* (2011) performed a numerical simulation of the flow over a NACA0014 airfoil undergoing heaving and pitching motions. They found out that compared to the sinusoidal motions, the use of non-sinusoidal flapping motions allowed to improve the power extraction by 17% and the efficiency was improved by 15%. These results were also confirmed by (Xiao *et al.* 2012), who highlighted the fact that the nature of the pitching motion had a great influence on the power extraction performance of a flapping NACA0012 airfoil. Moreover, according to their results, the trapezoidal pitching motion improved the output power coefficient by 63% and the efficiency by 50%. Later, Lu *et al.* (2014) showed that an appropriate combination of non-sinusoidal heaving and pitching motions enhanced the energy extraction performance. Their numerical results have shown that the use of a square-like pitching trajectory combined with a toothed-like heaving trajectory, improves the output power coefficient by 87.5% compared to the sinusoidal trajectory.

Through the literature, aerodynamic and biological studies have shown that the performances of a flapping airfoil were also influenced by the foil shape. Some researchers were inspired by the biological flexible structure of the insect wings or fish fins to improve the power extraction performance. Among them, [Liu *et al.* \(2013\)](#) that carried out numerical simulation to investigate the effect of the flexibility on the power extraction performance of a flapping airfoil. Their results showed that using a flexible leading edge and/or a flexible trailing edge enhanced the power efficiency by 7.68% compared to the rigid wing. [Wu *et al.* \(2015\)](#) used attached flexible tail at the trailing edge to improve the power extraction of a semi-actuated flapping foil. The tail used was either rigid or passively deformable. They observed that the net power extraction efficiency of the flexible tail was higher than the rigid tail. Then, [Xie *et al.* \(2016\)](#) have numerically explored the effects of the gurney flap on the energy extraction performances of flapping airfoil. Their results revealed that, compared to a standard NACA0012 flapping airfoil, the use of optimized gurney flap led to a significant increase of the output power coefficient and the energy extraction efficiency.

Furthermore, modifications of the foil shape were also considered as means to improve the aerodynamic or the hydrodynamic characteristics of the flapping foil. [Usuh *et al.* \(2012\)](#) investigated the flow over a non-profiled (rectangular) flapping plate to examine the foil shape effect. Their numerical results showed that the use of a non-profiled plate improves the efficiency of power extraction by 5.35% as compared to NACA0012 profile in similar flapping conditions. They also found out that the thickness of rectangular section has a negligible effect on the power extraction performances. [Le *et al.* \(2013\)](#) carried out two-dimensional Navier-Stokes simulation to explore the morphological effect of bio-inspired scallop-shell. Their results showed that an optimal foil shape with corrugation and camber improved the efficiency by 6% compared to a standard NACA0012 foil. Recently, [Wang *et al.* \(2016\)](#), using an in-house UCFD (Unified Computational Fluid Dynamics) software, investigated the influence of the geometrical parameters including the maximum foil thickness, maximum camber, position of maximum thickness and position of maximum camber of a NACA 4 digits and 6 digits series foils on the energy extraction performance of flapping hydrofoil. Their results showed that the performances of the flapping foils increased with the increase of the maximum thickness until an optimum value, and then, decreased with further increase of maximum thickness. For the symmetrical hydrofoils with same maximum thickness, but with different positions, it was observed that the hydrofoil performances increased first, and then decreased when shifting the maximum thickness position from the leading edge to the trailing edge. The best energy extraction performances were obtained at 1/4 chord maximum thickness position. While, for non-symmetric foils with same thickness, the use of larger camber

reduce the energy extraction performances.

Having a step on the foil surface is another technique of geometrical modification that was used to improve the aerodynamic characteristics. The idea of using a stepped airfoil was first introduced by [Fogleman and Kline \(1972\)](#). Thereafter, several experimental and numerical studies were performed to evaluate the performances of this new designed foil. [Fertis \(1994\)](#) investigated through a series of experimental tests on a three-dimensional wing model. His results showed considerable improvement of the aerodynamic characteristics in terms of lift, drag, lift-to-drag ratio and stall angle of attack. [Boroomand and Hosseini-verdi \(2009\)](#) investigated numerically the aerodynamic characteristics of stepped NACA-2412 airfoil. They concluded that, making a step on the lower surface, near the trailing edge, delays the static stall and slightly increases the drag force. However, the lift to drag ratio was improved. More recently, [Kamyab *et al.* \(2016\)](#) and [Kamyab and Ghassemi \(2017\)](#) performed a numerical and experimental study of the flow around a stepped NACA0012 airfoil with step on upper surface and both upper and lower surfaces at low Reynolds number. Their results also showed that using step in airfoil leads to a delay in the static stall, an increase in the lift coefficient and a decrease in the Strouhal number that resulted in lower noise generation on the airfoil.

To the authors' knowledge, no research has been yet carried out to study the aerodynamic characteristics of a stepped flapping foil. Therefore, in this paper, a new stepped foil shape is proposed to improve the energy extraction performance of a flapping foil. The new foil shape considered in this study is composed from two symmetric foils, which are the NACA0012 profile at the leading edge and the NACA006 profile at the trailing edge (Fig. 1). Leading and trailing segments are combined at a well-defined zone along the chord length. The assembly of the two half-profiles creates a step on both top and bottom surfaces. These steps change the vortex activities near the foil surface during the flapping cycle.

The main goal of the present work is to demonstrate the capability of the combined foil shape (or stepped foil) to passively control the vortex activities near the foil surface during the flapping cycle in order to improve the hydrodynamic characteristics. The study was performed using the computational fluid dynamics (CFD) code ANSYS Fluent 16. The solution of governing equations was obtained at low Reynolds number ($Re=10\ 000$). The imposed heaving and pitching motions are achieved by using dynamic mesh techniques and User Define Function (UDF).

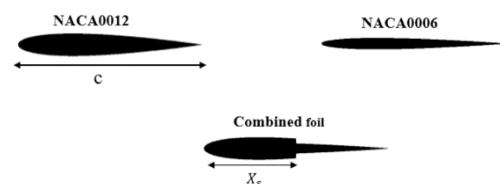


Fig. 1. Geometry of the studied foils.

2. PROBLEM DESCRIPTION AND NUMERICAL METHOD

2.1 Kinematics of Flapping Foil

The flapping motion is produced by the combination of two periodic motions : heaving $h(t)$ and pitching $\theta(t)$. The heaving is a vertical translational motion while the pitching is a rotational motion of the foil around an axis located on the chord line at the distance X_P from the leading edge. Figure 2 illustrates the flapping motions and their kinematic parameters.

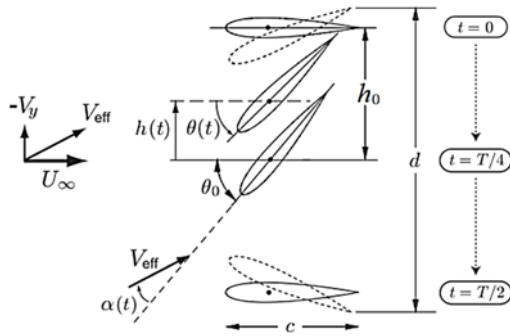


Fig. 2. Main kinematic parameters of a flapping foil (Kinsey and Dumas 2014).

In this study, the imposed flapping motion is performed according to sinusoidal trajectories described mathematically by the following equations:

$$h(t) = h_0 \sin(\omega t + \phi) \quad (1)$$

$$\theta(t) = \theta_0 \sin(\omega t) \quad (2)$$

Where h_0 and θ_0 are the maximum amplitudes of heaving and pitching respectively. $\omega = 2\pi f$ is the angular frequency and ϕ is the phase angle between heaving and pitching. c is the chord length.

In addition, two dimensionless parameters are used to describe the kinematics of flapping foil, Strouhal number defined by Anderson *et al.* (1998) as follows : ($St = 2fch_0/U_\infty$) and Reynolds numbers based on the chord length ($Re = \rho c U_\infty / \mu$). ρ and μ are the fluid density and dynamic viscosity respectively.

The flapping motion in a horizontal flow produces an effective angle of attack α_{eff} and an effective upstream velocity V_{eff} . These parameters have significant effects on the hydrodynamic force of flapping foil (Kinsey and Dumas 2008). These are given by the following relations:

$$\alpha_{eff}(t) = \arctan\left(\frac{-1}{U_\infty} \frac{dh(t)}{dt}\right) - \theta(t) \quad (3)$$

$$V_{eff}(t) = \sqrt{U_\infty^2 + \left(\frac{dh(t)}{dt}\right)^2} \quad (4)$$

Where U_∞ is the freestream velocity and $dh(t)/dt$ is the heaving velocity.

Flapping foil can operate in two different modes, propulsion or power extraction, depending on the foil orientation at the effective angle of attack. The operating mode is identified by the dimensionless parameter χ (feathering parameter) defined by Anderson *et al.* (1998) as:

$$\chi = \theta_0 / \arctan\left(\frac{h_0 \omega}{U_\infty}\right) \quad (5)$$

When the foil is in propulsion mode, then $\chi < 1$ while in power extraction mode $\chi > 1$. For a given flapping frequency the operating mode depends primarily on the maximum amplitude of the pitching.

2.2 Definition of Power Extraction and Efficiency

The instantaneous power $P(t)$ extracted from the flow by the flapping foil is written as:

$$P(t) = P_h(t) + P_\theta(t) \quad (6)$$

where $P_h(t)$ is the power extracted by the heaving motion and $P_\theta(t)$ is the power extracted by the pitching motion:

$$P_h(t) = F_y(t) \frac{dh(t)}{dt} \quad (7)$$

And

$$P_\theta(t) = M_z(t) \frac{d\theta(t)}{dt} \quad (8)$$

$F_y(t)$ is the instantaneous vertical force and $M_z(t)$ is the instantaneous moment.

The instantaneous power coefficient C_{OP} is defined as:

$$C_{OP} = C_{OPh} + C_{OP\theta} = \frac{P_h(t) + P_\theta(t)}{0.5\rho c U_\infty^3} \quad (9)$$

This parameter is also written as:

$$C_{OP} = \frac{1}{U_\infty} \left(C_L \frac{dh(t)}{dt} + C_M c \frac{d\theta(t)}{dt} \right) \quad (10)$$

where C_L and C_M are the lift and the moment coefficients, respectively.

The mean power extracted (\bar{P}) and the mean power coefficient (\bar{C}_{OP}) over one flapping cycle of period T are calculated as:

$$\bar{P} = \frac{1}{T} \left(\int_t^{t+T} F_y(t) \frac{dh(t)}{dt} dt + \int_t^{t+T} M_z(t) \frac{d\theta(t)}{dt} dt \right) \quad (11)$$

$$\bar{C}_{OP} = \frac{1}{T} \int_t^{t+T} C_{OP}(t) dt \quad (12)$$

The efficiency of power extraction η is defined as the ratio of the mean extracted power and the power available in the swept fluid area :

$$\eta = \frac{\bar{P}}{P_a} = \bar{C}_{OP} \frac{c}{d} \quad (13)$$

where $(P_a = 1/2 \rho U_\infty^3 d)$ is the maximum available power in the flow, d is the maximum vertical displacement of the trailing edge.

2.3 Algorithms and Models

A numerical simulation was performed using the commercial CFD code ANSYS-Fluent 16 based on the finite volume method. A segregated pressure solver was used to solve the 2-D unsteady and incompressible Navier-Stokes equations on a moving grid. The pressure-velocity coupling was achieved by means of the SIMPLEC algorithm, and a Green-Gauss Node-Based approach was used for gradient evaluation. The discretization of pressure and momentum terms are based on the second order scheme and the third order MUSCL scheme respectively. For temporal discretization, the first order implicit scheme was used. All simulations are carried out at a Reynolds number, $Re = 10^4$.

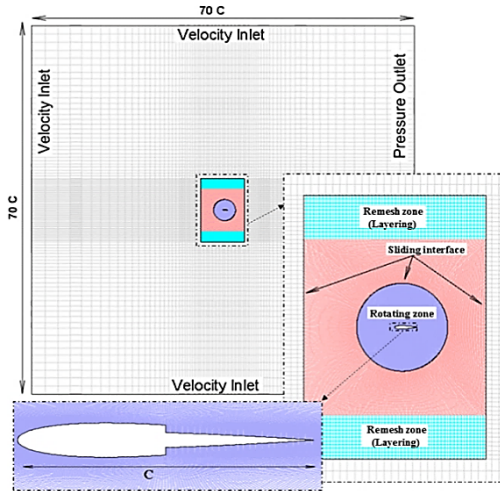


Fig. 3. Mesh and Boundary conditions.

Structured cells are built for the meshing of all the computational domain that is divided in three zones: a rotating zone, a layered zone and a fixed zone. The rotating zone is circular and contains the foil. It is meshed with a very fine quadrilateral structured grid in order to capture the physical gradients near the wall accurately. The first grid point in the normal direction of the foil surface is set to $4.610^{-5}c$. This first zone ensures the pitching and the heaving of the foil simultaneously or separately. The second zone (translating zone) translates vertically to ensure the heaving motion. In the fixed zone, the mesh is less dense. The connections

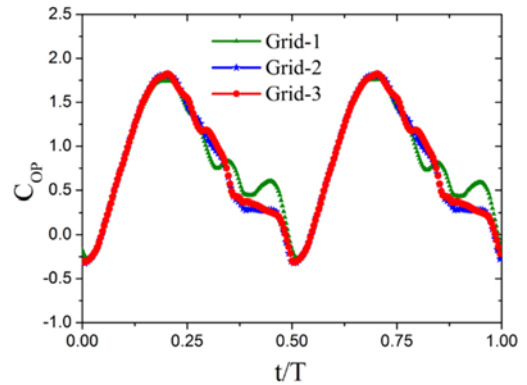
between the different zones are managed with non-conformal interfaces.

Dynamic meshes and sliding meshes techniques are applied to ensure the pitching and heaving motions. The dynamic mesh motion was controlled by a User Defined Function (UDF).

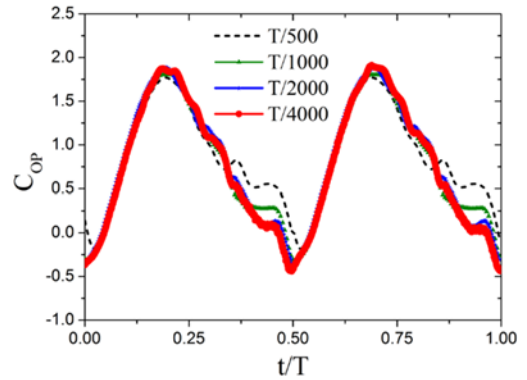
2.5 Sensitivity Study and Validation

A sensitivity study is carried out to ensure the independence of the numerical solution on the grid and the time step (Δt).

First, the dependence of the spatial resolution on the numerical results was explored. The mesh quality has a noticeable effect on the numerical simulation results. This is more accentuated when the grid is deformable. So, we consider three grids of different densities: Grid-1 is composed of 47 880 cells (with 200 points on the airfoil surface), Grid-2 consists of 87 760 cells (with 400 points on the airfoil surface) and Grid-3 consists of 254 400 cells (with 800 points on the airfoil surface). These simulations are performed at $Re = 10000$, $St = 0.3$, $h_0 = 1$, $\phi = 90^\circ$, $\theta = 73.30^\circ$ and $\Delta t = T/1000$. Figure 4-a represents the instantaneous power coefficient over one flapping cycle computed with the three grids. According to this figure, the difference between the results obtained with Grid-2 and Grid-3 is negligible. Thus, Grid-2 is sufficiently refined to obtain reliable results. Then, it will be applied for all the calculations that follow.



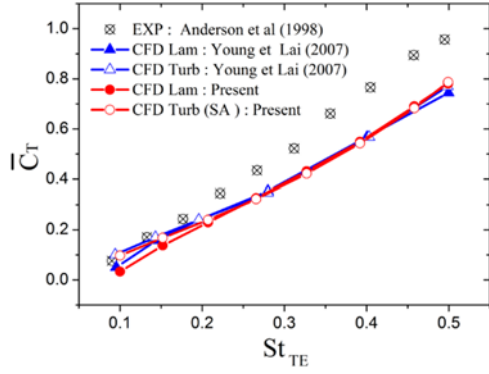
(a)



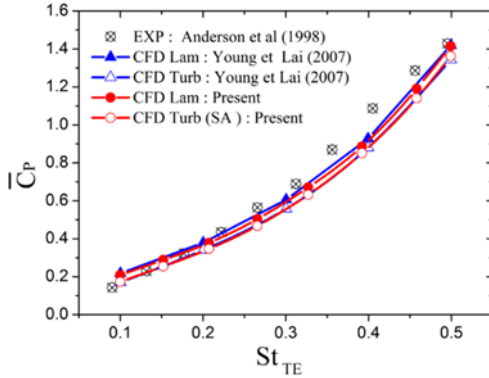
(b)

Fig. 4. Instantaneous power coefficient over one flapping cycle computed using (a) different grids and (b) different time steps .

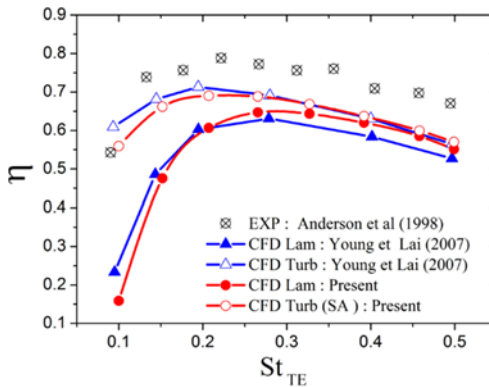
Then, an analysis of the temporal resolution is performed to find the appropriate time step size. Figure 4-b displays the variation of the instantaneous power coefficient during one flapping cycle computed with four time steps $T/500$, $T/1000$, $T/2000$ and $T/4000$. As it is seen in Fig. 4-b the numerical solutions are effectively independent on the time step size if the time step used is equal or higher than $T/2000$. Therefore, the time-step $\Delta t = T/2000$ is set to all following simulations.



(a) \bar{C}_T



(b) \bar{C}_P



(c) η

Fig. 5. Comparison of (a) mean thrust coefficient, (b) mean power coefficient, and (c) propulsive efficiency versus St at $h_0 = 0.75$, $\alpha_0 = 15^\circ$, $\phi = 90^\circ$ and 410^4 .

In addition to the sensitivity study, a validation study is performed by comparing the computed results to experimental and numerical results published by Anderson *et al.* (1998) and Young and Lai (2007).

These numerical simulations are carried out for a 2D incompressible flow over NACA0012 foil undergoing combined heaving and pitching motions at low Reynolds number 410^4 . The motion parameters are those of the experimental study performed by Anderson *et al.* (1998): $h_0 = 0.75$, $\alpha_0 = 15^\circ$ and $\phi = 90^\circ$. The flapping frequency has been varied in the range $[0.268 - 1.34]$ to ensure a Strouhal number between 0.1 and 0.5. In this validation study, the Strouhal number is based on the trailing edge excursion, $St_{TE} = 2fd/U_\infty$, as in (Anderson *et al.* 1998). The computed mean thrust coefficient (\bar{C}_T), input power coefficient

(\bar{C}_P) and (b) the propulsive efficiency (η) are compared to the experimental results of Anderson *et al.* (1998) and to those obtained numerically by Young and Lai (2007). Figure 5 shows that the numerical values are slightly lower than the experimental results but with the same tendency. Differences between numerical and experimental results were also found in other CFD studies (Guglielmini and Blondeaux 2004, Xiao and Liao 2010, Karbasian *et al.* 2015). As reported by Guglielmini and Blondeaux (2004), this difference can be due to the inaccurate procedure used by Anderson *et al.* (1998) during the experimental measurement of forces and to the 2D approximation used in the numerical simulations. However, a good agreement is obtained between our results and the numerical results published by Young and Lai (2007).

3. RESULTS AND DISCUSSIONS

According to Kinsey and Dumas (2008), Zhu and Peng (2009), Ashraf *et al.* (2011) and Xiao *et al.* (2012), the best energy extraction performances of a flapping foil are obtained when using the following kinematic parameters: $St = 0.2$ to 0.4 , $\theta_0 = 60^\circ$ to 90° , $X_p = 0.25c$ to $0.5c$, $h_0 = c$ and $\phi = 90^\circ$. These values give favourable conditions to the formation and shedding process of the leading edge vortex, and also ensure a good synchronization between the lift force and the heaving velocity. Hence, in this study, all simulations are carried out using the following kinematic parameters: $h_0 = 1c$, $\phi = 90^\circ$, $X_p = 0.333$, $\alpha_{eff} \in [10^\circ - 35^\circ]$, $St \in [0.15 - 0.45]$ and $Re = 10000$. All results obtained for the combined foil shape are compared to those of NACA0012 and NACA0006 base foils.

3.1 Effect of the Connecting Area

In this section, the effect of the connecting area location (X_s/c) is investigated. For this purpose, four locations are considered $X_s/c = 0.2, 0.33, 0.5$ and 0.8 respectively. These computations are performed with $St = 0.4$ and $\alpha_{eff} = 30^\circ$.

The variations of \bar{C}_{OPh} , $\bar{C}_{OP\theta}$, \bar{C}_{OP} and η are displayed in Fig. 6 as a function of X_s/c . Figs. 6-a to 6-c show that the connecting area location has a great effect on the flapping foil performances and that when a combined foil shape is used, the energy extraction performances of the flapping foil are improved. Moreover, Figs. 6-a and 6-b show that \bar{C}_{OPh} has a positive contribution to the power extraction while the contribution from the pitching

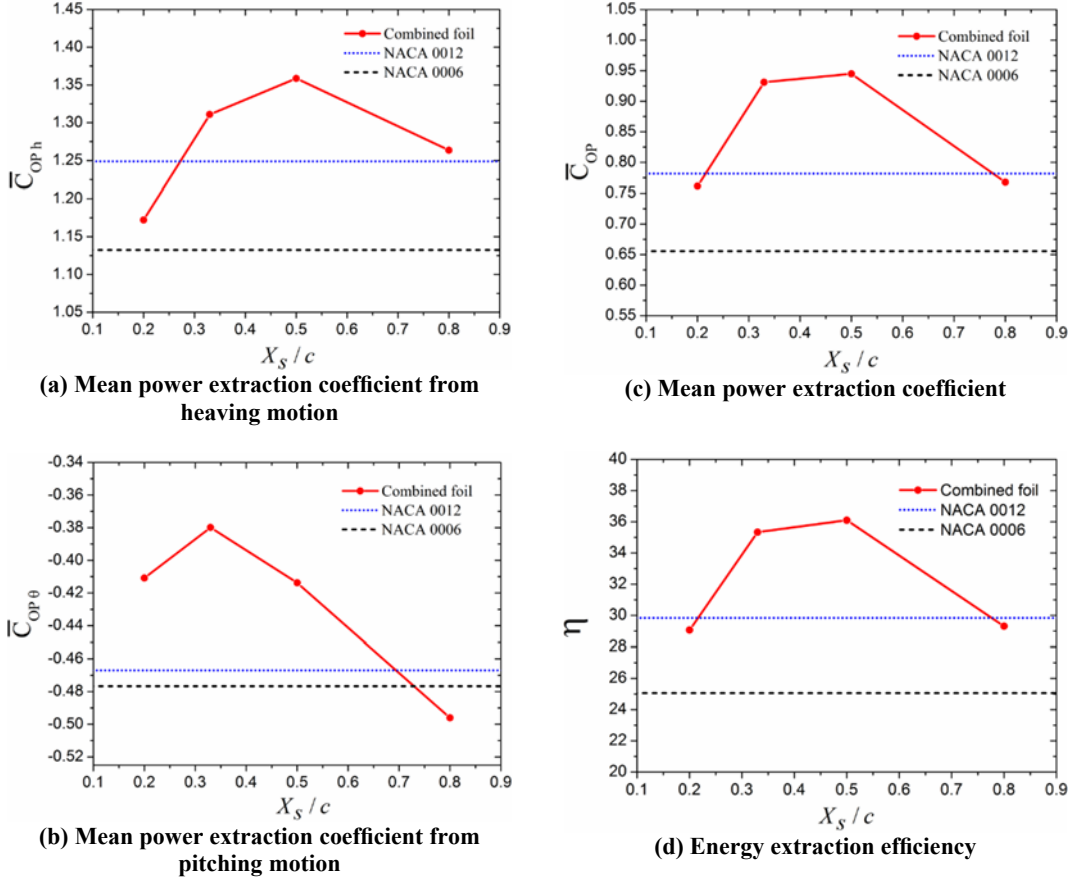


Fig. 6. Comparison of \bar{C}_{OPh} , $\bar{C}_{OP\theta}$, \bar{C}_{OP} and η versus X_s/c at $St = 0.4$ and $\alpha_{eff} = 30^\circ$.

motion is negative ($\bar{C}_{OPh} \leq 0$). Thus the pitching motion consumes energy. Therefore, the total energy extracted by the flapping foil is dominated by the energy extracted through the heaving motion. It is also found that with the use of an optimal combined foil shape, the energy consumed by the pitching motion is reduced and hereby, the energy extraction performances of the flapping foil are improved. From Figs. 6-c and 6-d, it can be seen that the best performances are obtained when the connecting area is located at the middle of the chord ($X_s/c = 0.5$). In this case, the power extraction efficiency is improved by 17.32% and 30.60% compared to the NACA0012 and NACA0006 baseline foils respectively (Fig. 6-d). Moving the connection area (X_s/c) toward the leading edge or the trailing edge considerably decreases the energy extraction performance of the combined flapping foil.

Figure 7 shows the temporal evolution of C_{OPh} , $C_{OP\theta}$ and C_{OP} over one flapping cycle for the combined foil with $X_s/c = 0.5$ and the baseline foils at $St=0.4$ and $\alpha_{eff} = 30^\circ$. It is quite clear that the combined foil shape has a significant effect on the power extraction coefficient. As shown in Fig. 7-a, the instantaneous C_{OPh} generated by the combined foil is higher than that of NACA0012 and NACA0006 baseline foils during the flapping time when $t = 0.1T$ to $t = 0.3T$ and $t = 0.6T$ to $t = 0.8T$,

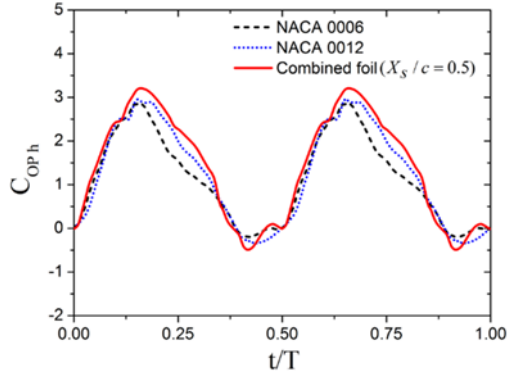
while in Fig. 7-b the $C_{OP\theta}$ is higher when $t = 0.4T$ to $0.5T$ and $t = 0.9T$ to $t = T$. As a result, the overall power extraction coefficient (C_{OP}) is increased (Fig. 7-c).

The energy extraction performances of a flapping foil are directly related to the variation of C_L and C_M coefficients. As seen in Eq. (10), the power extracted from the heaving motion is determined by the product of the lift coefficient and the heaving velocity. While the power extracted from the pitching motion is determined by the product of the moment coefficient and the angular velocity.

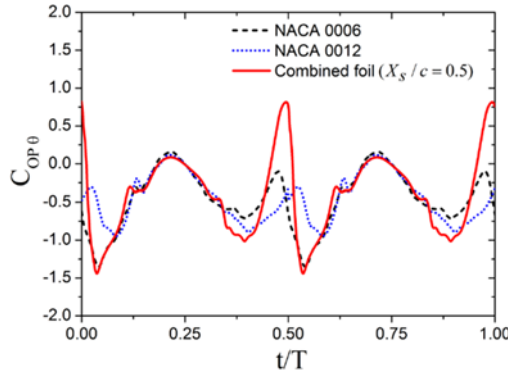
The time variation of C_L , $((dh/dt)/U_\infty)$, C_M , $((cd\theta/dt)/U_\infty)$ and C_D over a flapping cycle are drawn in Fig. 8 for the NACA0006 and NACA0012 baseline foils, and the combined foil with $X_s/c = 0.5$.

Figure 8-a shows that, for the three foil shapes, the lift force and the heaving motion have the same sign during the whole of flapping cycle. This improves the energy extracted from the heaving motion. It is also shown that the instantaneous lift produced by the combined foil is higher than those of the baseline foils. However, the moment coefficient and the angular velocity have opposite signs (Fig. 8-c). Therefore, their effect on the mean power is negative. It is also observed that the moment generated by the combined foil is higher than that

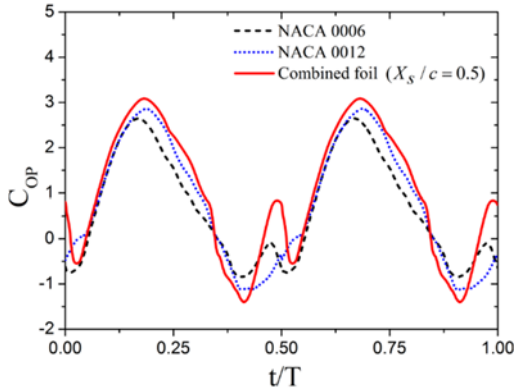
generated by the baseline foils. These results are in agreement with those of DeLaurier and Harris (1974) that the use of a stepped airfoil increases the moment coefficient.



(a) power extraction coefficient from heaving motion



(b) power extraction coefficient from pitching motion

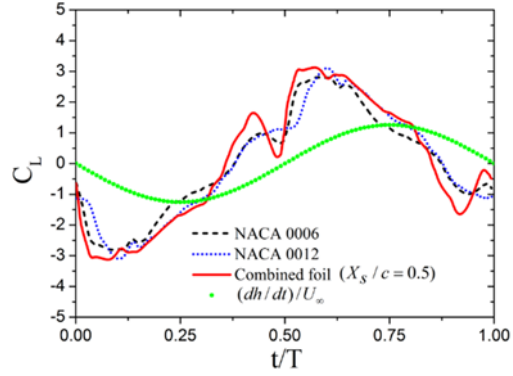


(c) total power extraction coefficient

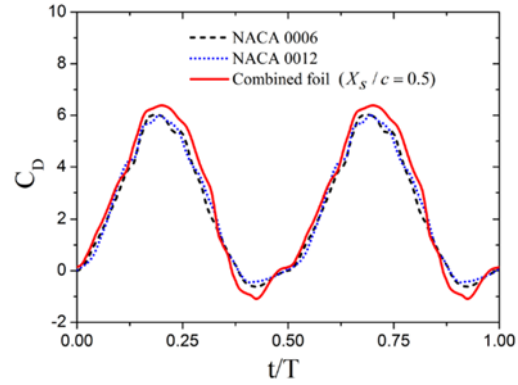
Fig. 7. Time evolution of (a) C_{Oph} , (b) C_{Op0} , and (c) C_{OP} over one flapping cycle at $St = 0.4$ and $\alpha_{eff} = 30^\circ$.

Figure 8-b presents the temporal variation of the drag coefficients (C_D) over one flapping cycle. As it can be observed that the drag coefficient of the combined foil increases slightly compared to the baseline foils. These results are consistent with those of (Fertis 1994, Boroomand and Hosseini-verdi 2009, Kamyab *et al.* 2016 and Kamyab and Ghassemi 2017) in which it was concluded that the stepped foil can improve the lift coefficient and the

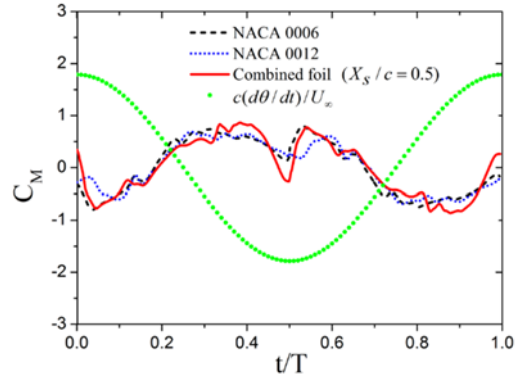
lift to drag ratio with a little increase of the drag coefficient. Xie *et al.* (2016) noted that in the field of energy extraction by flapping foil, more attention is focused to the lift improvement, while the drag penalty can be offset by the structural amelioration.



(a)



(b)



(c)

Fig. 8. Time evolution of (a) lift coefficient (C_L) and heaving velocity ($(dh/dt)/U_\infty$), (b) drag coefficient (C_D), (c) Moment coefficient (C_M) and pitching velocity ($(cd0/dt)/U_\infty$) over one flapping cycle at $St = 0.4$ and $\alpha_{eff} = 30^\circ$.

Therefore, the main differences in forces and moment distribution are attributed to the change in flow structure around the foil caused by the geometric modification. Detailed analysis of the flow fields around the flapping foil helps to clarify how the combined foil shape (or stepped foil) improves the energy extraction performance.

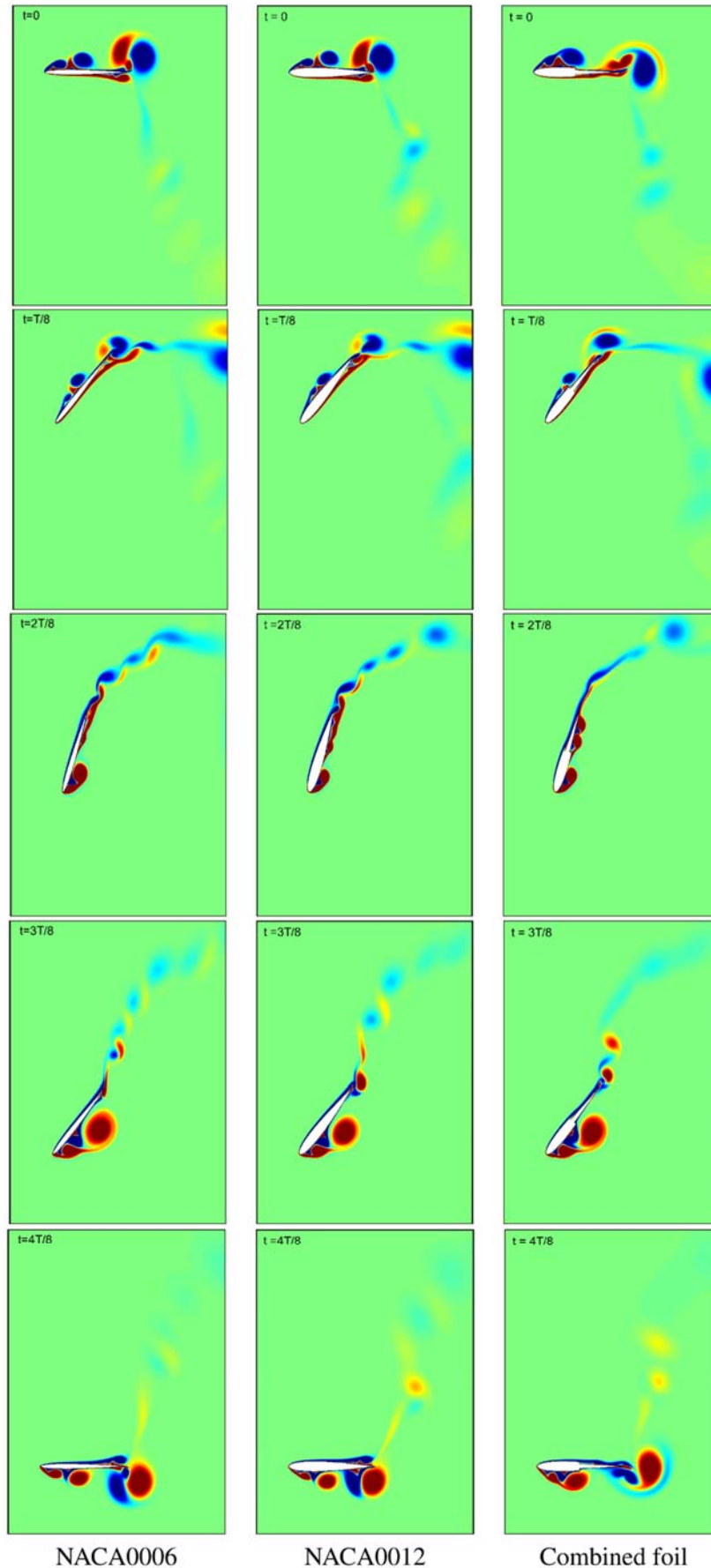


Fig. 9. Instantaneous vorticity contours around the NACA0006, NACA0012 and the combined foil during one half flapping cycle at $St = 0.4$ and $\alpha_{eff} = 30^\circ$.

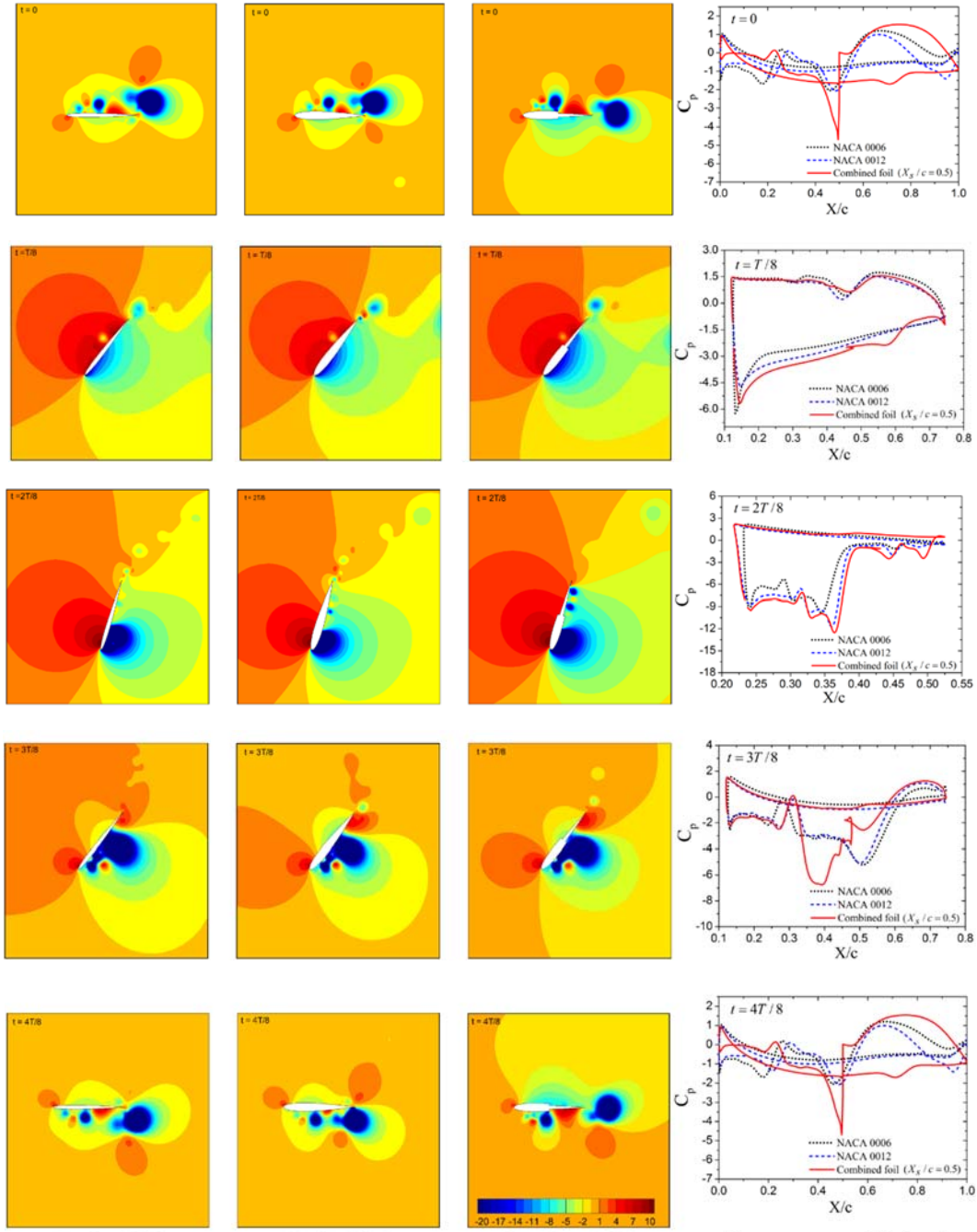
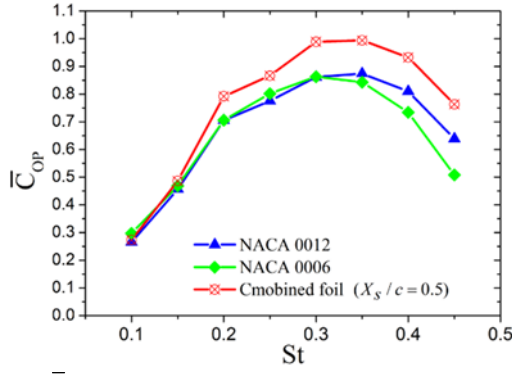


Fig. 10. Instantaneous pressure contours and pressure coefficients around the NACA0006, NACA0012 and the combined foil during one half flapping cycle at $St = 0.4$ and $\alpha_{eff} = 30^\circ$.

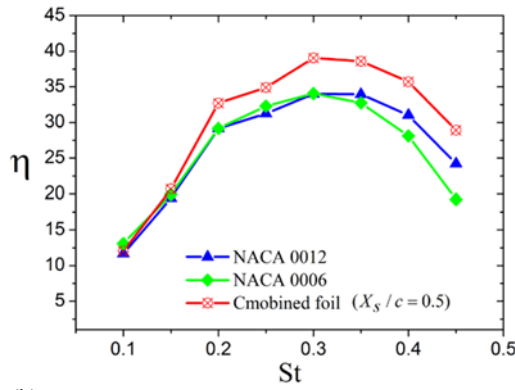
Figure 8, shows that C_L , C_D and C_M coefficients are symmetric for both downstroke and upstroke phases. Therefore, only results of the first half period are presented and analyzed. Figures 9 and 10 show respectively the instantaneous vorticity contours and the instantaneous pressure contours and the pressure coefficient around the NACA0006, NACA0012 and the combined foil during one half flapping cycle at $St = 0.4$ and $\alpha_{eff} = 30^\circ$. It was found that the step has a little effect on the LEV formation process, but this effect is significant on the vortex activity near the trailing edge.

At $t = 0$, the flapping foils are at their maximum positive position of the heaving motion and begin to downstroke. It can be seen that the leading edge vortex, formed during the previous period, interacts with the other part of the foil before shedding in the wake.

During the flapping time $t = 0$ to $t = 2T/8$, the flapping foils are far from the highest position of the heaving motion and simultaneously pitching in the counterclockwise direction. The LEV continues to grow on the top surfaces and convects downstream from the foils.



(a) \bar{C}_{OP}



(b) η

Fig. 11. Comparison of (a) mean power extraction coefficient and (b) energy extraction efficiency versus St at $\alpha_{eff} = 30^\circ$.

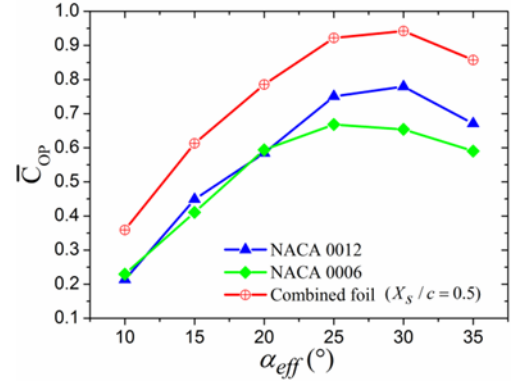
At $t = 2T/8$, the pitching angle and the heaving velocity reach their maximum values. At the same time, a leading edge vortex (LEV) starts to form and a low-pressure region is observed near the leading edge on the

bottom surface of the foils. However, at time $t = 2T/8$, the combined foil generates a high pressure zone on the top surface and a low pressure zone on the bottom surface. Furthermore, a large lifting force is generated comparatively to the base foils. Consequently, C_{OPh} obtained with the combined foil at this time is higher.

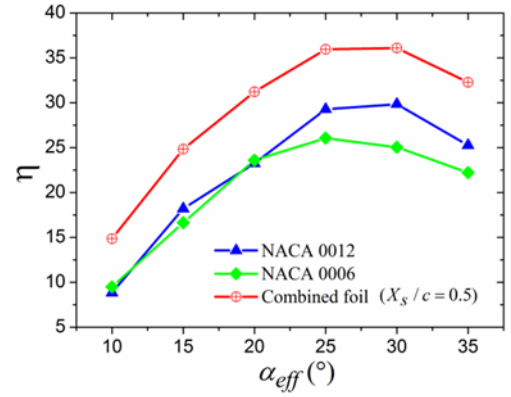
During the flapping time from $2T/8$ to $4T/8$, the LEV is convected downstream along the foil surfaces and the low-pressure region is shifted downstream.

At $t = 4T/8$, the LEV interacts with the trailing edge, and the force caused by the low pressure center is applied near the trailing edge. At this time, the steps on the combined foil surface accelerate the LEV shedding. Consequently, the vortex configuration near the trailing edge is modified. Synchronization between the moment coefficient and the angular velocity is improved and $C_{OP\theta}$ generated by the pitching motion is significantly increased.

Thereafter, this process is repeated symmetrically in the second half period (upstroke phase).



(a) \bar{C}_{OP}



(b) η

Fig. 12. Comparison of (a) mean power extraction coefficient and (b) energy extraction efficiency versus α_{eff} at $St = 0.4$.

3.2 Effect of Strouhal Number

The variations of the mean power coefficient and the energy extraction efficiency as a function of the Strouhal number are shown in Fig. 11. It was found that both C_{OP} and η increase with St until a maximum values at $St = 0.3$ and then decrease. The further increase of St causes an important decrease of the energy extraction performances. These results are explained by the good synchronization between the lift forces and the heaving motion at St close to 0.3. Conversely, for higher or lower values of St , the synchronization is weaker. Moreover, at high St number the energy consumed by the pitching motion increases rapidly compared to the energy extracted from the heaving motion. It was also noticed that at low Strouhal number ($St < 0.2$) the foil shape has a negligible effect on the energy extraction performance. But for $St > 0.2$, significant improvements of C_{OP} and η are found with the combined foil.

3.3 Effect of the Angle of Attack (AOA)

The variation of the mean output power coefficient and the energy extraction efficiency according to the effective angle of attack is shown in Fig. 12. The Strouhal number was set to $St = 0.4$. For the three flapping foils, when the effective angle of attack increases, both C_{OP} and η increase first and then decrease. The optimal angles of attack are

between 25° and 35°. Moreover, the mean output power coefficient and the energy extraction efficiency were found to be higher for the combined foil compared to NACA0006 and NACA0012 baseline foils for all angles of attack. This confirms that the combined foil shape can improve the energy extraction performance of a flapping foil.

4. CONCLUSION

In this work, the energy extraction performance of a flapping foil device with a combined foil shape was investigated. The 2D unsteady and incompressible Navier-Stokes equations governing the flow over the foil undergoing prescribed heaving and pitching motions at low Reynolds number ($Re = 10\,000$) were solved by a finite volume method, using the software ANSYS Fluent 16. The considered foil was designed by combining a NACA0012 leading edge and a NACA006 trailing edge. Then, the NACA0006 and NACA0012 were used as baseline foils. For this work, a validation study was carried out, and the results were in good agreement with the published data. Thus, the main results obtained in this study are summarized below:

- 1) The use of the combined foil improves the energy extraction performances.
- 2) The connecting area between the leading edge and the trailing edge (X_s/c) has a significant effect on the energy extraction performance. Indeed, the best performance was observed at $X_s/c = 0.5$, where the energy extraction efficiency was improved by 30.60% and 17.32% in comparison to NACA0006 and NACA0012, respectively.
- 3) For the kinematic parameters considered in this study, it is also observed that the combined flapping foil extracts the energy mainly by the heaving motion, while the contribution of the pitching motion is negative.
- 4) Moreover, the flow fields over a flapping cycle show that the vortex and the pressure distribution near the trailing edge of the combined foil are changed. This is due to the stepped surface of the foil. Consequently, the lift force and the moment coefficient are enhanced, thus, the energy extraction performance is improved.

Accordingly, the combined foil shape may provide an effective means to improve the energy extraction performance.

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